

Higgs inflation and Quantum Gravity

Michał Artymowski

University of Warsaw

June 29, 2012

Collaboration with A. Dappor and T. Pawłowski
Erice 2012

Outline

Outline

- ▶ Motivation

Outline

- ▶ Motivation
- ▶ Inflation and Higgs inflation

Outline

- ▶ Motivation
- ▶ Inflation and Higgs inflation
- ▶ Loop Quantum Cosmology

Outline

- ▶ Motivation
- ▶ Inflation and Higgs inflation
- ▶ Loop Quantum Cosmology
- ▶ LQC and non-minimal coupling to gravity

Outline

- ▶ Motivation
- ▶ Inflation and Higgs inflation
- ▶ Loop Quantum Cosmology
- ▶ LQC and non-minimal coupling to gravity
- ▶ Summary

Motivation

Motivation

- ▶ Matching the data

Motivation

- ▶ Matching the data
- ▶ Theoretical advantages

Motivation

- ▶ Matching the data
- ▶ Theoretical advantages
- ▶ New physics at gravitational level

Motivation

- ▶ Matching the data
- ▶ Theoretical advantages
- ▶ New physics at gravitational level
- ▶ Strong quantum gravity effects

Inflation

Isotropic, homogeneous universe \Rightarrow FRW metric

$$g_{\mu\nu} = \text{diag} (1, -a^2, -a^2, -a^2) , \quad (1)$$

where $a = a(t)$ is a scale factor.

Inflation

Isotropic, homogeneous universe \Rightarrow FRW metric

$$g_{\mu\nu} = \text{diag} (1, -a^2, -a^2, -a^2) , \quad (1)$$

where $a = a(t)$ is a scale factor. Inflation appears, when $\ddot{a} > 0$, but in particular we are interest in

$$H = \frac{\dot{a}}{a} \simeq \text{const} \Rightarrow a \sim e^{Ht}$$

Inflation

Isotropic, homogeneous universe \Rightarrow FRW metric

$$g_{\mu\nu} = \text{diag} (1, -a^2, -a^2, -a^2) , \quad (1)$$

where $a = a(t)$ is a scale factor. Inflation appears, when $\ddot{a} > 0$, but in particular we are interest in

$$H = \frac{\dot{a}}{a} \simeq \text{const} \Rightarrow a \sim e^{Ht}$$

Actually it's enough if the slow-roll parameters are small

$$\epsilon = -\frac{\dot{H}}{H^2} , \quad \eta = \frac{\ddot{H}}{\dot{H}H} , \quad \epsilon, |\eta| \ll 1 \quad (2)$$

Inflation

Isotropic, homogeneous universe \Rightarrow FRW metric

$$g_{\mu\nu} = \text{diag} (1, -a^2, -a^2, -a^2) , \quad (1)$$

where $a = a(t)$ is a scale factor. Inflation appears, when $\ddot{a} > 0$, but in particular we are interest in

$$H = \frac{\dot{a}}{a} \simeq \text{const} \Rightarrow a \sim e^{Ht}$$

Actually it's enough if the slow-roll parameters are small

$$\epsilon = -\frac{\dot{H}}{H^2} , \quad \eta = \frac{\ddot{H}}{\dot{H}H} , \quad \epsilon, |\eta| \ll 1 \quad (2)$$

Example: massive scalar field domination

Inflation

Isotropic, homogeneous universe \Rightarrow FRW metric

$$g_{\mu\nu} = \text{diag} (1, -a^2, -a^2, -a^2) , \quad (1)$$

where $a = a(t)$ is a scale factor. Inflation appears, when $\ddot{a} > 0$, but in particular we are interest in

$$H = \frac{\dot{a}}{a} \simeq \text{const} \Rightarrow a \sim e^{Ht}$$

Actually it's enough if the slow-roll parameters are small

$$\epsilon = -\frac{\dot{H}}{H^2} , \quad \eta = \frac{\ddot{H}}{\dot{H}H} , \quad \epsilon, |\eta| \ll 1 \quad (2)$$

Example: massive scalar field domination

Problem: $m \sim 10^{13} \text{ GeV}$ - way beyond SM physics

Inflation from the non-minimal coupling

The action in Jordan frame is of the form of

$$S = \int d^4x \sqrt{-g} \left[-U(\phi)R + \frac{1}{2}(\partial_\mu\phi)(\partial^\mu\phi) - V(\phi) \right], \quad (3)$$

where $U = \frac{1}{2} + f(\phi)$ and $f(\phi)R$ is a non-minimal coupling to gravity. EOMs from Lagrange equations.

Inflation from the non-minimal coupling

The action in Jordan frame is of the form of

$$S = \int d^4x \sqrt{-g} \left[-U(\phi)R + \frac{1}{2}(\partial_\mu\phi)(\partial^\mu\phi) - V(\phi) \right], \quad (3)$$

where $U = \frac{1}{2} + f(\phi)$ and $f(\phi)R$ is a non-minimal coupling to gravity. EOMs from Lagrange equations. Let us assume, that $U = \frac{1}{2}(1 + \xi\phi^2)$ and $V(\phi) = \frac{1}{4}\lambda\phi^4 - \frac{1}{2}m^2\phi^2$.

For $\xi\phi^2 \gg 1$, $\lambda \gg \xi m^2$ and $\dot{\phi}^2 \ll V(\phi)$ one obtains

$$\epsilon \sim \eta \sim \frac{1}{\xi\phi^2} \ll 1, \quad \xi = 47000\sqrt{\lambda}$$

Inflation from the non-minimal coupling

The action in Jordan frame is of the form of

$$S = \int d^4x \sqrt{-g} \left[-U(\phi)R + \frac{1}{2}(\partial_\mu\phi)(\partial^\mu\phi) - V(\phi) \right], \quad (3)$$

where $U = \frac{1}{2} + f(\phi)$ and $f(\phi)R$ is a non-minimal coupling to gravity. EOMs from Lagrange equations. Let us assume, that $U = \frac{1}{2}(1 + \xi\phi^2)$ and $V(\phi) = \frac{1}{4}\lambda\phi^4 - \frac{1}{2}m^2\phi^2$.

For $\xi\phi^2 \gg 1$, $\lambda \gg \xi m^2$ and $\dot{\phi}^2 \ll V(\phi)$ one obtains

$$\epsilon \sim \eta \sim \frac{1}{\xi\phi^2} \ll 1, \quad \xi = 47000\sqrt{\lambda}$$

One obtains accelerated expansion of space-time!

Inflation ends for $\phi \sim \xi^{-1/2} \ll M_{pl}$

Loop Quantum Gravity corrections

One can describe gravity by the Hamiltonian and Hamilton equations. Let us define Ashtekar canonical variables c, p and their Hamiltonian by

$$p = a^2, \quad c = \gamma \dot{a}, \quad \mathcal{H} = -\frac{3}{\gamma^2} c^2 \sqrt{p} + \mathcal{H}_{\text{matt}}. \quad (4)$$

Classical equations of motion can be obtained from

$$\dot{c} = \{c, \mathcal{H}\}, \quad \dot{p} = \{p, \mathcal{H}\}, \quad \{c, p\} = \frac{\gamma}{3} \quad (5)$$

Loop Quantum Gravity corrections

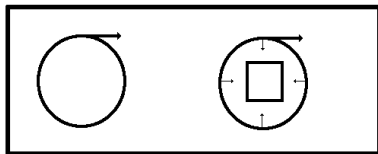
One can describe gravity by the Hamiltonian and Hamilton equations. Let us define Ashtekar canonical variables c, p and their Hamiltonian by

$$p = a^2, \quad c = \gamma \dot{a}, \quad \mathcal{H} = -\frac{3}{\gamma^2} c^2 \sqrt{p} + \mathcal{H}_{\text{matt}}. \quad (4)$$

Classical equations of motion can be obtained from

$$\dot{c} = \{c, \mathcal{H}\}, \quad \dot{p} = \{p, \mathcal{H}\}, \quad \{c, p\} = \frac{\gamma}{3} \quad (5)$$

The loop correction



Loop Quantum Gravity corrections

Assumption: only LQC holonomy correction is important. Then

$$\pi_\phi \rightarrow \pi_\phi, \quad \rho \rightarrow \rho, \quad c \rightarrow \frac{1}{\bar{\mu}} \sin(\bar{\mu}c), \quad (6)$$

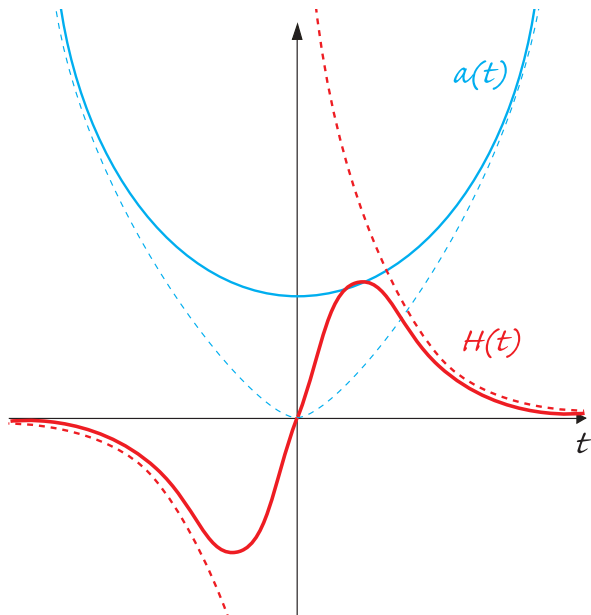
where $\bar{\mu} = \sqrt{\Delta/\rho}$ and Δ is a minimal value of area operator. This gives Friedmann equations of the form of

$$3H^2 = \rho \left(1 - \frac{\rho}{\rho_{cr}}\right), \quad 2\dot{H} = -(\rho + P) \left(1 - 2\frac{\rho}{\rho_{cr}}\right), \quad (7)$$

where $\rho_{cr} = \frac{3}{\Delta\gamma^2}$ is a critical (maximal) energy density.

This leads to the Big Bounce for $\rho = \rho_{cr}$

LQC in FRW Universe



Einstein frame

Let us redefine $g_{\mu\nu}$ and $\phi(t)$, so new scalar field would be canonical and minimally coupled

$$\tilde{g}_{\mu\nu} = 2U(\phi)g_{\mu\nu}, \quad \left(\frac{d\tilde{\phi}}{d\phi}\right)^2 = \frac{1}{2} \frac{U + 3U'^2}{U^2}, \quad \tilde{V}(\tilde{\phi}) = \frac{1}{4} \frac{V(\phi)}{U^2(\phi)}. \quad (8)$$

Einstein frame

Let us redefine $g_{\mu\nu}$ and $\phi(t)$, so new scalar field would be canonical and minimally coupled

$$\tilde{g}_{\mu\nu} = 2U(\phi)g_{\mu\nu}, \quad \left(\frac{d\tilde{\phi}}{d\phi}\right)^2 = \frac{1}{2} \frac{U + 3U'^2}{U^2}, \quad \tilde{V}(\tilde{\phi}) = \frac{1}{4} \frac{V(\phi)}{U^2(\phi)}. \quad (8)$$

New action $\tilde{S}[\tilde{g}_{\mu\nu}, \tilde{\phi}] = S[g_{\mu\nu}, \phi]$ has a minimal coupling $\tilde{U} = \frac{1}{2}$.

Einstein frame

Let us redefine $g_{\mu\nu}$ and $\phi(t)$, so new scalar field would be canonical and minimally coupled

$$\tilde{g}_{\mu\nu} = 2U(\phi)g_{\mu\nu}, \quad \left(\frac{d\tilde{\phi}}{d\phi}\right)^2 = \frac{1}{2} \frac{U + 3U'^2}{U^2}, \quad \tilde{V}(\tilde{\phi}) = \frac{1}{4} \frac{V(\phi)}{U^2(\phi)}. \quad (8)$$

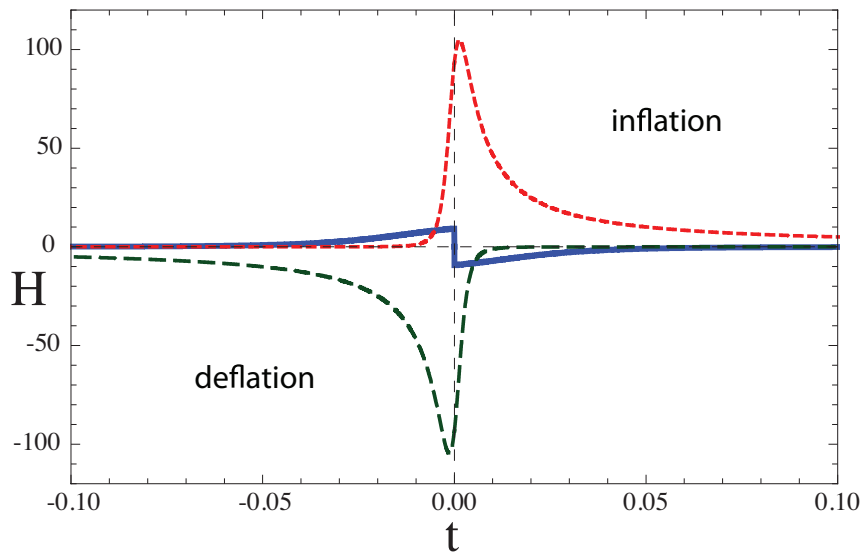
New action $\tilde{S}[\tilde{g}_{\mu\nu}, \tilde{\phi}] = S[g_{\mu\nu}, \phi]$ has a minimal coupling $\tilde{U} = \frac{1}{2}$.

Together with LQC holonomy correction one finds

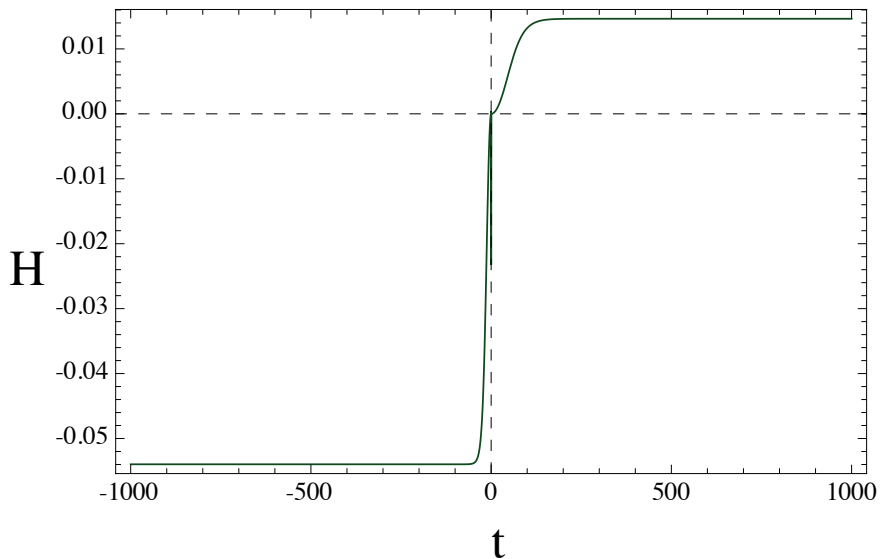
$$3\tilde{H}^2 = \tilde{\rho} \left(1 - \frac{\tilde{\rho}}{\rho_{cr}}\right), \quad 2\frac{d\tilde{H}}{d\tilde{t}} = -(\tilde{\rho} + \tilde{P}) \left(1 - 2\frac{\tilde{\rho}}{\rho_{cr}}\right). \quad (9)$$

Big Bounce for $\tilde{\rho} = \rho_{cr}$!

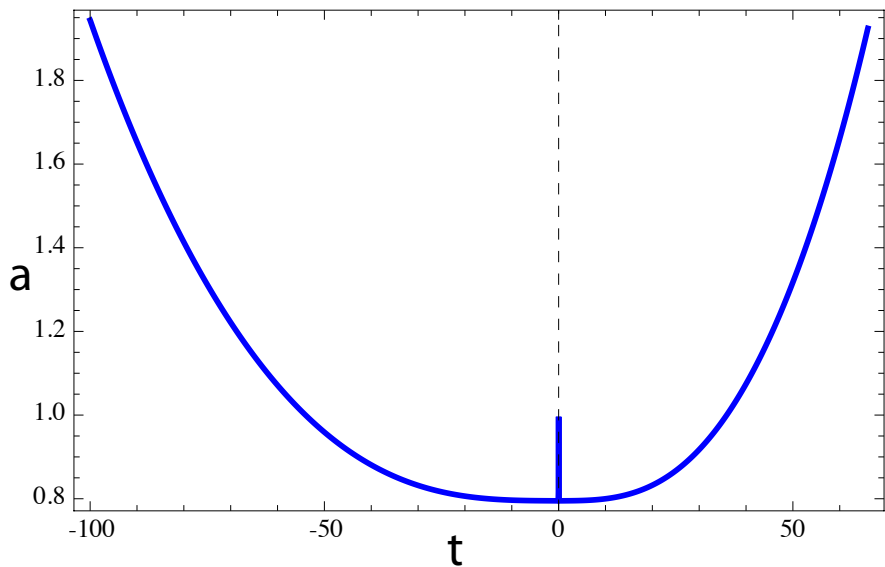
Evolution of Hubble parameters



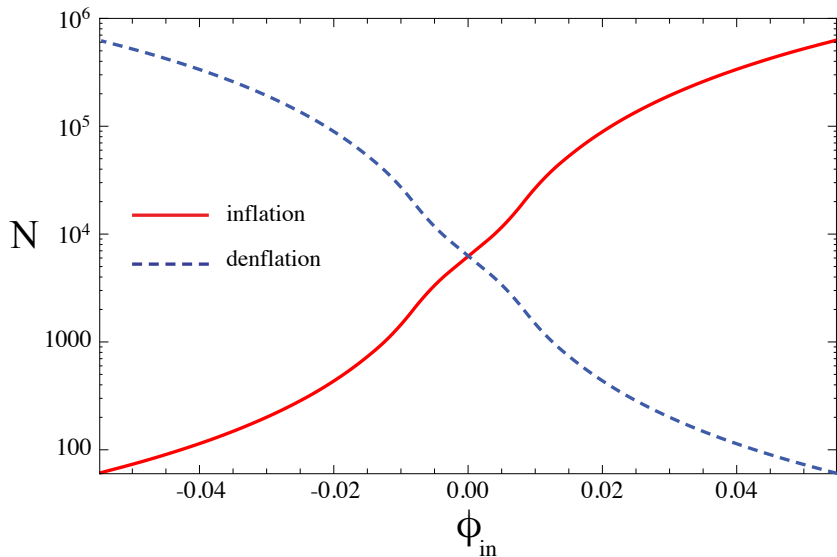
Evolution of the Hubble parameter



Evolution of a scale factor



e-folds



Summary

Summary

- ▶ Mass of standard inflaton is beyond any typical scale of SM

Summary

- ▶ Mass of standard inflaton is beyond any typical scale of SM
- ▶ Solution: non-minimal coupling of Higgs field

Summary

- ▶ Mass of standard inflaton is beyond any typical scale of SM
- ▶ Solution: non-minimal coupling of Higgs field
- ▶ Very small ϵ, η - long and stable inflation

Summary

- ▶ Mass of standard inflaton is beyond any typical scale of SM
- ▶ Solution: non-minimal coupling of Higgs field
- ▶ Very small ϵ, η - long and stable inflation
- ▶ Strong quantum gravity effects: multiple bounces and anadiabatic evolution

Summary

- ▶ Mass of standard inflaton is beyond any typical scale of SM
- ▶ Solution: non-minimal coupling of Higgs field
- ▶ Very small ϵ, η - long and stable inflation
- ▶ Strong quantum gravity effects: multiple bounces and anadiabatic evolution
- ▶ Natural initial conditions for inflation