

# Tractor beams, Cosmic strings, and related phenomena



Péter Forgács<sup>1,2</sup>, Árpád Lukács<sup>1</sup>, Tomasz Romańczukiewicz<sup>3</sup>

<sup>1</sup> MTA Wigner RCP Institute for Particle and Nuclear Physics, H1525 Budapest, POB 49

<sup>2</sup> LMPT, CNRS-UMR 6083, Université de Tours, Parc de Grandmont, 37200 Tours, France

<sup>3</sup> Institute of Physics, Jagiellonian University, Reymonta 4, 30-059 Cracow, Poland

lukacs.arpad@wigner.mta.hu, www.rmki.kfki.hu/~arpi, arXiv:1303.3237



## 1. Introduction

Forces acting on a scatterer

- multichannel problem
- energy and momentum conservation

**Tractor beam or Negative radiation pressure (NRP):** force on the scatterer acts in the opposite direction as the wave travels

Other approaches

- **gain media** Mizraki, Fainman, [1]
- **structured beam** scattered dominantly into the direction behind the scatterer, Sukhov, Dogariu [2]
- **Negative index materials** force opposite to wave number, but not to Poynting vector [3]
- **Higher harmonics** due to **nonlinearities**, Romańczukiewicz; Forgács, Lukács and Romańczukiewicz [4]

## 2. One dimensional examples

Two channels,  $u$  and  $d$ , for channel  $u$ :

$$F_u/A_u^2 = k_u(1 + |R_{uu}|^2 - |T_{uu}|^2) + k_d(|R_{du}|^2 - |T_{du}|^2), \quad (1)$$

- $R_{ij}$  reflection coefficient, incoming wave  $j$  into outgoing mode  $i$
- $T_{ij}$  transmission coefficient similarly
- $k_u, k_d$ : wave numbers

Eq. (1) follows immediately by taking into account that a plane wave  $A \exp(ikx)/\sqrt{k}$ , carries momentum  $k|A|^2$ . We note that the energy flux of such a wave is  $\omega|A|^2$ .

Energy conservation:  $\sum_i (|R_{ij}|^2 + |T_{ij}|^2) = 1$ .

### A macroscopic example

- Also applicable to electromagnetic wave propagation
- Material with  $n_x > n_y$
- Propagation in  $z$  direction, modes:  $E_x, E_y$  **polarizations**
- Scatterer: same material rotated by  $45^\circ$  around  $z$  axis
- Incoming wave of energy flow  $10^3 \text{ W/cm}^2 = 10^7 \text{ W/m}^2$  in  $x$  polarization radiation pressure:  $-1.05 \text{ Pa}$
- For total reflection :  $2.45 \text{ Pa}$

### A birefringent platelet

- The same system for light scattering
- outer medium: ordered liquid crystal,  $n_x > n_y$
- Scatterer: birefringent, axes  $45^\circ$  rotated wrt. axes of outer medium
- using data of 5CB, wavelength  $5893 \text{ \AA}$ ,  $0.1 \text{ mm}$  thick scatterer
- NRP:  $x$  polarization  $-3.04 \times 10^{-12} \text{ Pa m/V}|E_0|^2$   
 $y$  polarization,  $2.40 \times 10^{-12} \text{ Pa m/V}|E_0|^2$
- few percents of light pressure (and transparent)
- light pressure:  $\sim 10 \mu\text{Pa}$  for sunlight

## 3. Force on the scatterer in 2D

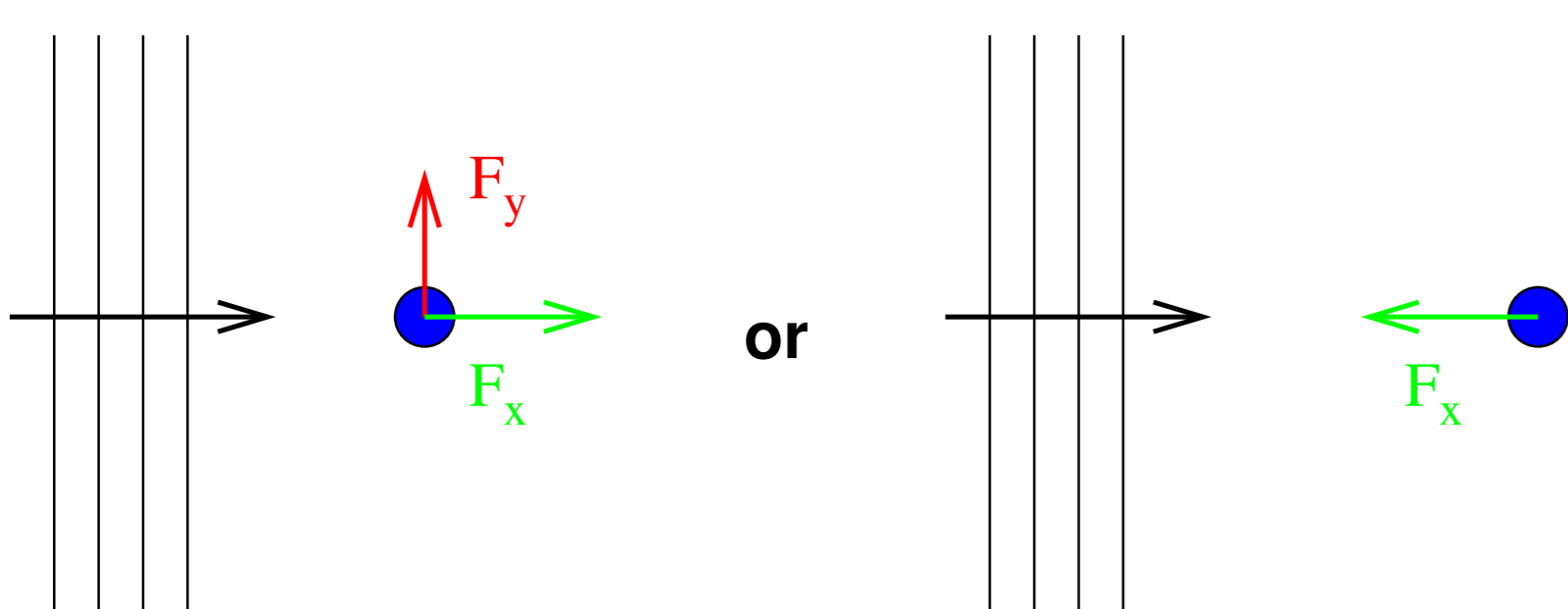


Figure 1: Force components

Force acting on the scatterer: **momentum flow**

$$\mathbf{F} = - \lim_{R \rightarrow \infty} \int_{-\pi}^{\pi} \mathbf{T}_e R d\vartheta, \quad (2)$$

assumptions: **free propagation** for  $R \rightarrow \infty$ . Inserting partial waves:

$$F = -4 \sum_{\ell} \left\{ A^{\dagger} S_{\ell+1}^{\dagger} K S_{\ell} A - A^{\dagger} K A \right\} \quad (3)$$

where  $F = F_x + iF_y$ ,  $K = \text{diag}(k_1, k_2)$ : **wave numbers** of modes,  $A = (A_1, A_2, \dots)^T$ .

Also applicable for **Aharonov-Bohm scattering** [5]:

- $S_{\ell} = \exp(2i\delta_{\ell})$ ,  $\delta_{\ell} = \pi(\ell - \nu)/2$ ,  $\nu^2 = (\ell - A_0)^2$ , where  $A_0$  is **no. of flux quanta**

- reproduces known results:  $F_x = -2k [\cos(2\pi A_0) - 1] |A|^2$ ,  $F_y = 2k \sin(2\pi A_0) |A|^2$

## 4. Scattering off cosmic strings

Cosmic strings[6]

- String core: Higgs field zero
- $\Phi \sim v \exp(i\vartheta)$  ( $v$ : VEV,  $\vartheta$ : angle)
- gauge covariant derivative  $\rightarrow 0$  for  $r \rightarrow \infty$ :  $\mathbf{A} \propto \mathbf{e}_{\vartheta}/r$  (Aharonov-Bohm-like asymptotics)
- GUT string: vector potential inside string corresponds to broken gauge generator (X boson)

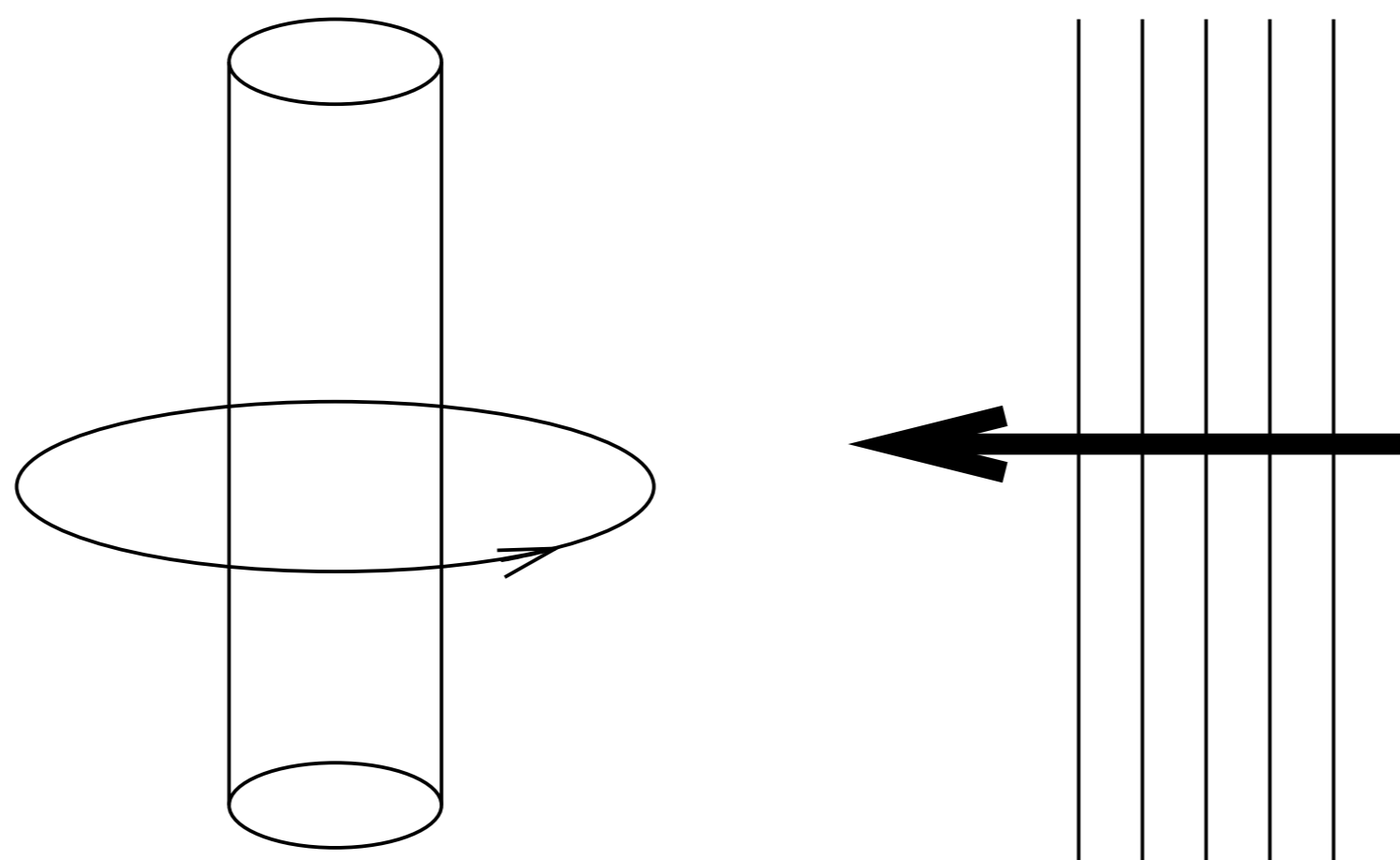


Figure 2: Scattering off a vortex line/cosmic string

Early Universe: friction dominated era of string network **Approximate force**: Aharonov-Bohm scattering [6]

- $F = - \sum_i 4n_i v (1 - \exp(2\pi i \nu_i \Phi))$   
where  $\nu_i$  is the charge (coupling strength between particle species  $i$  and the X-boson)
- incoming flux,  $|A_i|^2 k_i = n_i v$ ,  $n_i$  is the number density of PS  $i$

**String induces baryon number violating processes**[7]

$$B + \text{string} \rightarrow \ell + \text{string} \quad (4)$$

Simplified model: ignore spins, one  $U(1)$  gauge potential

$$(\nabla + i\mathbf{A} \frac{\sigma_2}{2})^2 \rho - K^2 \rho = 0, \quad \sigma_2 = \begin{pmatrix} & -i \\ i & \end{pmatrix}, \quad (5)$$

where

- $\mathbf{A} = \mathbf{e}_{\vartheta}/r$  gauge potential of the string
- $\rho = (u, d)$  wave function of two species,  $u$  heavy,  $d$  light
- fermionic boundary conditions  $\rho(r, \vartheta + 2\pi) = -\rho(r, \vartheta)$
- units:  $m_u = 2$ ,  $m_d = 1.5$ ,  $\hbar = c = 1$

Partial waves:

$$(u, d) = \sum_{\ell=-\infty}^{\infty} e^{i(\ell+\gamma)\vartheta} (u_{\ell}(r), d_{\ell}(r)), \quad (6)$$

with  $\gamma = 1/2$ . The radial functions satisfy

$$\begin{aligned} u_{\ell}'' + \frac{u_{\ell}'}{r} - \frac{\eta_u^2}{r^2} u_{\ell} + \frac{c}{r^2} d_{\ell} + k_u^2 u_{\ell} &= 0, \\ d_{\ell}'' + \frac{d_{\ell}'}{r} - \frac{\eta_d^2}{r^2} d_{\ell} + \frac{c^*}{r^2} u_{\ell} + k_d^2 d_{\ell} &= 0, \end{aligned} \quad (7)$$

where  $\eta_u^2 = \eta_d^2 = (\ell + 1/2)^2 + 1/4$  and  $c = i(\ell + 1/2)$ .

## 5. Force acting on the cosmic string

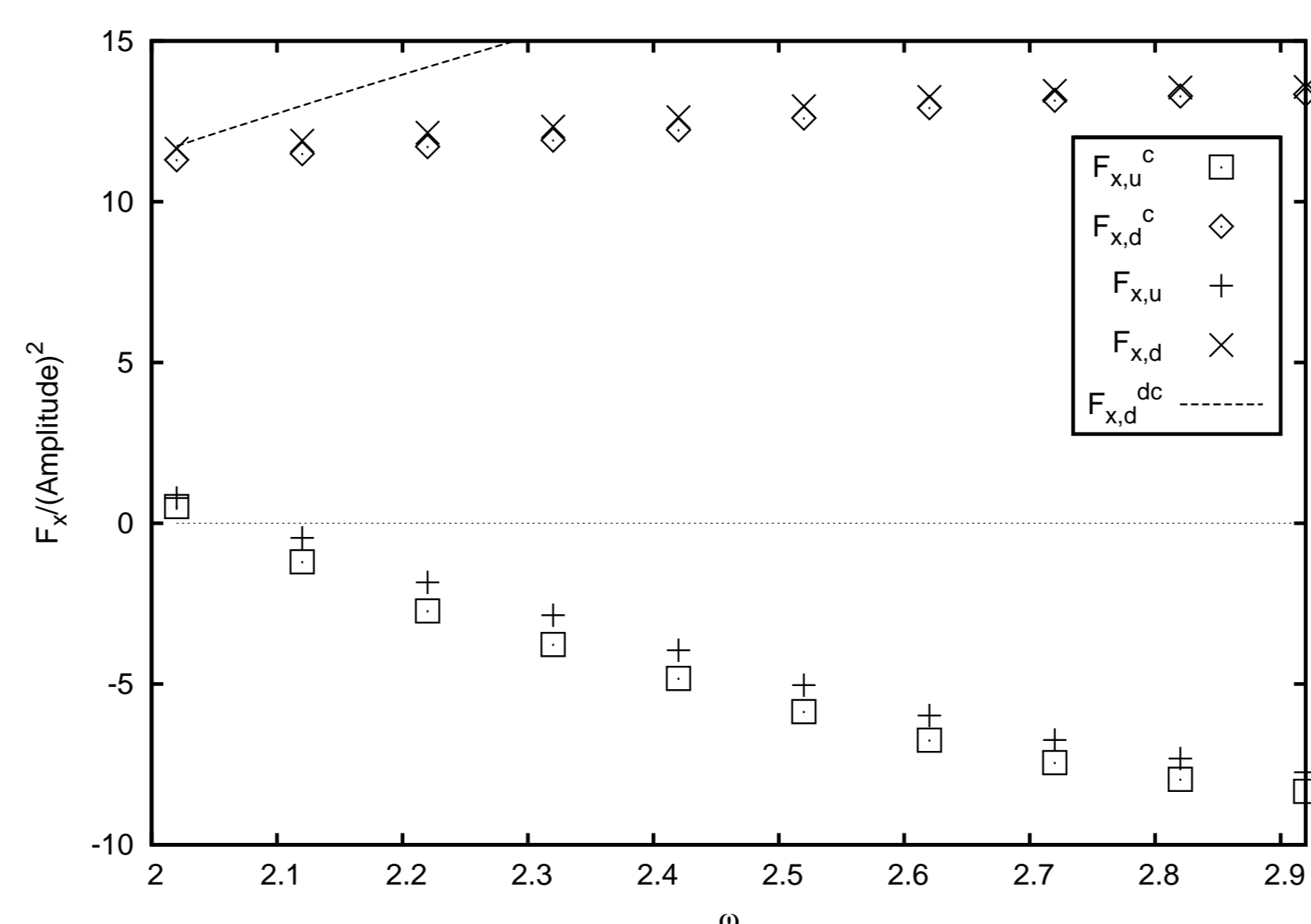


Figure 3: The longitudinal force component  $F_x$  as a function of the frequency  $\omega$ : the numerical result is compared to the decoupled, (dc) approximation, the latter only depicted for  $2 < \omega \leq 2.3$ . For the incoming heavy mode  $u$ , the radiation pressure becomes negative at  $\omega \approx 2.1557$ . The effect of the vortex core is also plotted: the longitudinal force  $F_x$  acting on a point vortex and on a vortex with a linear core (marked with 'c'). The existence of a vortex core enhances NRP.

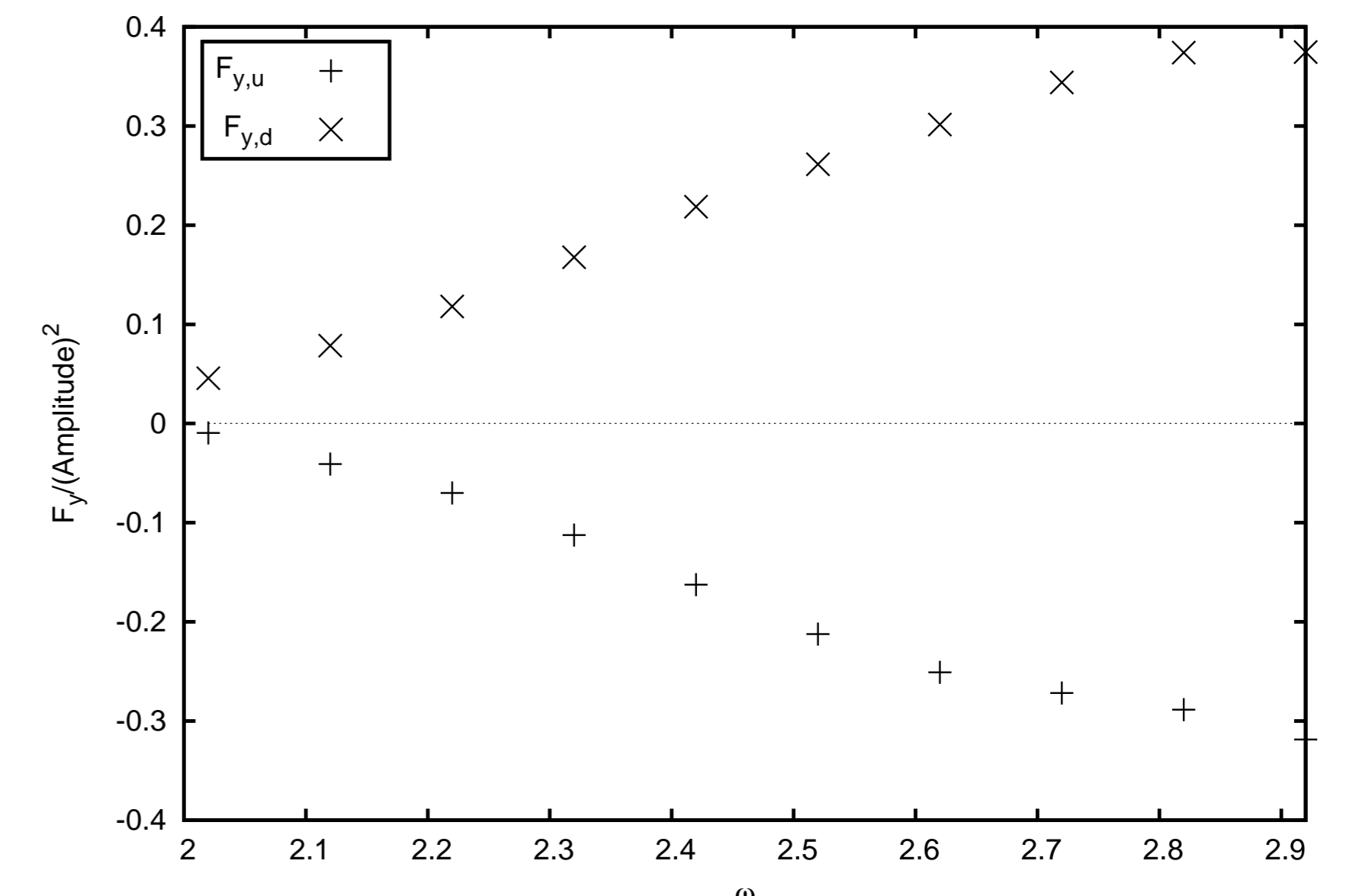


Figure 4: The transversal force component  $F_y$  as function of the frequency  $\omega$ .

- numerical solution of eqns. (7)
- AB-approximation: OK for light mode, **even the sign is wrong** for heavy one
- Effect is robust: vortex core does not influence it much
- At  $v = 0.655c$  (typical value at the end of friction era)

$$F_x^u = -6.09|A|^2, \quad F_x^d = 7.44|A|^2.$$

Total force: add up forces for each particle species

- amplitude:  $|A_i|^2 k_i \propto n_i v$
- scattering energy as  $m_i/\sqrt{1-v^2}$   
where  $v$  is the string velocity

### NRP on cosmic strings

- Baryon number violating scattering on cosmic strings
- Heavy baryonic modes accelerate the string instead of acting as a medium with friction
- Longer path
- More baryon asymmetry washed away

## 6. Further application

Planar scattering situations

- Reproduced known results for superfluid vortices (one channel)
- vortices in superconductors
- preliminary results: two-channel scattering in two-band superconductors (e.g.,  $\text{MgB}_2$ ), scattering between channels influences force greatly, but no tractor beam

## References

- [1] Mizraki, A. and Fainman, Y., *Opt. Lett.* **35** 3405 (2010).
- [2] Sukhov, S. and Dogariu, A. *Phys. Rev. Lett.* **107** 203602 (2011).
- [3] V.G. Veselago, *Sov. Phys. Usp.* **10** (4): 509–14 (1968).
- [4] P. Forgács, Á. Lukács, T. Romańczukiewicz, *Phys. Rev.* **D77**, 125012 (2008).
- [5] Y. Aharonov and D. Bohm, *Phys. Rev.* **115** (1959) 485. S. Olariu and I.I. Popescu, *Rev. Mod. Phys.* **57** 339-436 (1985); A.L. Shelankov, *Europhys. Lett.* **43** (6) pp. 623-628 (1998).
- [6] A. Vilenkin and E.P.S. Shellard, *Cosmic strings and other topological defects*, Cambridge University Press, Cambridge, 1994.
- [7] M.G. Ahlford, J. March-Russell, and F. Wilczek, *Nucl. Phys.* **B328** (1989) 140-158; W.B. Perkins, L. Perivolaropoulos, A.-C. Davis, R.H. Brandenberger and A. Matheson, *Nucl. Phys.* **B353** (1990) 237-270.