

Topological Invariant For Star-Triangle Relation

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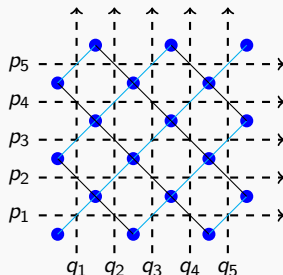
June, 2018

Introduction For Models in Statistical Mechanics

The aim of statistical mechanics is to predict the relations between the observable macroscopic properties of the system, given only a knowledge of the microscopic forces between the components.

- ▶ What do we want to understand?
- ▶ What do we want to calculate?

Two Dimensional Exactly Solved Models



- ▶ On each site i we place a 'spin' σ_i which takes some set of discrete or continuous values.
- ▶ For Ising Model:
 $\sigma_i = \pm 1$
 Periodic boundary condition $\sigma_{n+1} = \sigma_1$ and no external field.
 Spin interaction occurs only between neighbour spins.

Boltzmann Weights

The Boltzmann weight for the Ising model of the edge $\langle i, j \rangle$

$$\mathcal{W}(\sigma_i, \sigma_j) = \exp\{-\varepsilon(\sigma_i, \sigma_j)/k_B T\}$$

Boltzmann weights for two-dimensional models, are characterized according to the spin variables and “rapidity” variables ρ and q .

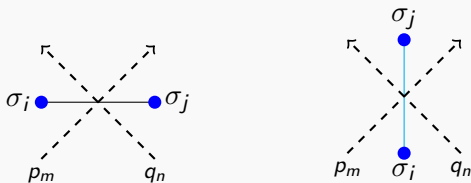


Figure: $\mathcal{W}_{\rho_m q_n}(\sigma_i, \sigma_j)$ (left) and $\overline{\mathcal{W}}_{\rho_m q_n}(\sigma_i, \sigma_j)$.

Transfer Matrices Commutation

Partition Function = $\text{Tr}(\mathcal{W}\overline{\mathcal{W}})$

Product $(\mathcal{W}\overline{\mathcal{W}})$ is not symmetric.

- ▶ Transfer matrices diagonalization

The Free energy of two dimensional Ising model in zero field was first obtained by Onsager in 1944. He diagonalized the transfer matrix by using matrix algebra.

Star-Triangle Relation

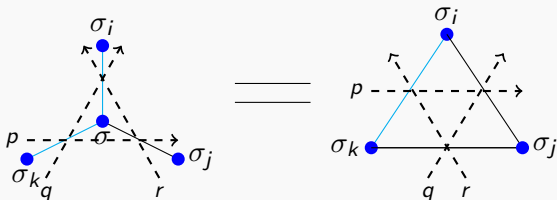


Figure: The star-triangle relation

Onsager(1971) noted that the star-triangle relation implies that their diagonal-to-diagonal transfer matrices commutation. Baxter(1978) has shown that 2D Ising model can be solved quite directly by star-triangle relation.

$$\sum_{\sigma} \mathcal{S}(\sigma) \overline{\mathcal{W}}_{qr}(\sigma, \sigma_i) \mathcal{W}_{pr}(\sigma_j, \sigma) \overline{\mathcal{W}}_{pq}(\sigma_k, \sigma)$$

$$= \mathcal{R}_{(pqr)} \mathcal{W}_{pq}(\sigma_i, \sigma_j) \overline{\mathcal{W}}_{pr}(\sigma_k, \sigma_i) \mathcal{W}_{qr}(\sigma_k, \sigma_j)$$

$$\begin{aligned} \sum_{\sigma} \mathcal{S}(\sigma) \overline{\mathcal{W}}_{qr}(\sigma, \sigma_i) \mathcal{W}_{pr}(\sigma_j, \sigma) \overline{\mathcal{W}}_{pq}(\sigma_k, \sigma) \\ = \mathcal{R}_{(pqr)} \mathcal{W}_{pq}(\sigma_i, \sigma_j) \overline{\mathcal{W}}_{pr}(\sigma_k, \sigma_i) \mathcal{W}_{qr}(\sigma_k, \sigma_j) \end{aligned}$$

The R_{pqr} is some factor independent of the spins.

- Spins can be interchanged.

There exist spin-independent quantities:

$$\prod_{\sigma_i} \mathcal{W}_{pq}(\sigma_i, \sigma_j) = \prod_{\sigma_i} \mathcal{W}_{pq}(\sigma_j, \sigma_i) = P_{pq}$$

$$\prod_{\sigma_i} \overline{\mathcal{W}}_{pq}(\sigma_i, \sigma_j) = \prod_{\sigma_i} \overline{\mathcal{W}}_{pq}(\sigma_j, \sigma_i) = \overline{P}_{pq}$$

for all values of the spin σ_i let Y_{pq} be the $N \times N$ matrix with entry $\mathcal{W}_{pq}(\sigma_i, \sigma_j)$ in position (σ_i, σ_j) , \overline{Y}_{pq} the matrix corresponding entry $\overline{\mathcal{W}}_{pq}(\sigma_i, \sigma_j)$.

Diagonal Matrices:

$$S = S(\sigma)\delta(\sigma_i, \sigma_j)$$

$$X_{pq|\sigma} = \mathcal{W}_{pq}(\sigma, \sigma_i)\delta(\sigma_i, \sigma_j), \quad X'_{pq|\sigma} = \mathcal{W}_{pq}(\sigma_i, \sigma)\delta(\sigma_i, \sigma_j)$$
$$\bar{X}_{pq|\sigma} = \bar{\mathcal{W}}_{pq}(\sigma, \sigma_i)\delta(\sigma_i, \sigma_j), \quad \bar{X}'_{pq|\sigma} = \bar{\mathcal{W}}_{pq}(\sigma_i, \sigma)\delta(\sigma_i, \sigma_j)$$

in position (σ_i, σ_j) . All diagonal matrices depend on the spin σ .

$$\begin{aligned} \sum_{\sigma} \mathcal{S}(\sigma) \overline{\mathcal{W}}_{qr}(\sigma, \sigma_i) \mathcal{W}_{pr}(\sigma_j, \sigma) \overline{\mathcal{W}}_{pq}(\sigma_k, \sigma) \\ = \mathcal{R}_{(pqr)} \mathcal{W}_{pq}(\sigma_i, \sigma_j) \overline{\mathcal{W}}_{pr}(\sigma_k, \sigma_i) \mathcal{W}_{qr}(\sigma_k, \sigma_j) \end{aligned}$$

Regard the σ_j as fixed and think of each side of the Star-Triangle Equation as the element of σ_i, σ_k of some matrix.

Matrix form of Star-Triangle relation:

$$Y_{qr} S X_{pr|\sigma_j} \overline{Y}_{pq} = \mathcal{R}_{pqr} X_{pq|\sigma_j} \overline{Y}_{pr} X_{qr|\sigma_j}$$

One can obtain \mathcal{R}_{pqr} by taking determinants.

$$\mathcal{R}_{pqr} = f_{qr} f_{pq} / f_{pr}$$

$$f_{pq} = P_{pq}^{-1} \det(S \overline{Y}_{pq})$$

The Invariant \mathcal{I}

Hold σ_i fixed and think them as elements (σ_j, σ_k)

$$Y_{pr} S X_{qr|\sigma_i} \bar{Y}_{pq} = \mathcal{R}_{pqr} X'_{pq|\sigma_i} Y_{qr} \bar{X}_{pr|\sigma_i}$$

$$\mathcal{R}_{pqr} = f_{pq} \bar{f}_{pr} / \bar{f}_{qr}$$

Hold σ_k fixed and regarding the relations as elements (σ_i, σ_j) of a matrix,

$$Y_{pr} S \bar{X}'_{pq|\sigma_k} \bar{Y}_{qr}^T = \mathcal{R}_{pqr} X'_{qr|\sigma_k} Y_{pq} \bar{X}'_{pr|\sigma_k}$$

$$\mathcal{R}_{pqr} = f_{qr} \bar{f}_{pr} / \bar{f}_{pq}$$

The relations of the \mathcal{R} 's are mutually consistent if and only if:

$$f_{qr}\bar{f}_{qr} = f_{pr}\bar{f}_{pr} = f_{pq}\bar{f}_{pq}$$

The only way this can happen is for the quantity

$$\mathcal{I} = N^{-1}f_{pq}\bar{f}_{pq}$$

to be independent of p and q!

Invariant \mathcal{I} for exactly solved models.

Ising model: $\mathcal{I}^2 = \sinh 2K \sinh 2\bar{K} = 1$ [R.J.Baxter.Exactly Solved Models in Statistical Mechanics. Academic, London, 1982.]

Ashkin-Teller model: $\mathcal{I} = 1$

Self-dual Potts model: $\mathcal{I} = 1$

Fateev-Zamolodchikov model : $\mathcal{I} = 1$

[Baxter 0108363]

Summary

- ▶ What we have learned?

Solutions to the star-triangle relation leads to two-dimensional exactly solvable models.

For all previously solved models \mathcal{I} is topological invariant, it does not depend on rapidity variables.

- ▶ Some Of The New Models

[Spiridonov 1011.3798]

[Kels 1504.07074]

[Gahramanov, Spiridonov 1505.00765]

Thank you for your attention!