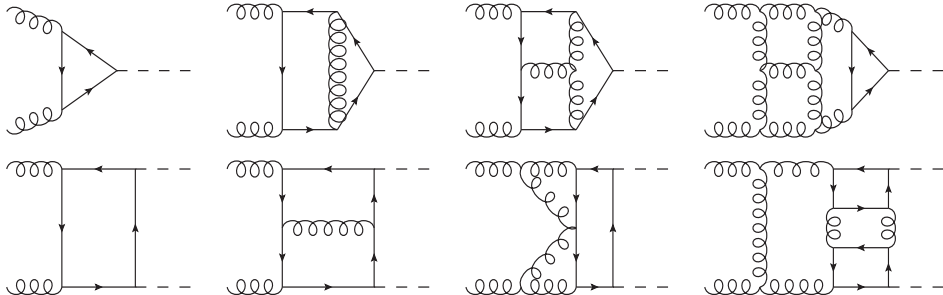


# Running, Decoupling & Higgs-Physics

Florian Herren | 15.06.2018

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INSTITUT FÜR THEORETISCHE TEILCHENPHYSIK



# Higgs-Boson production in Gluon-fusion

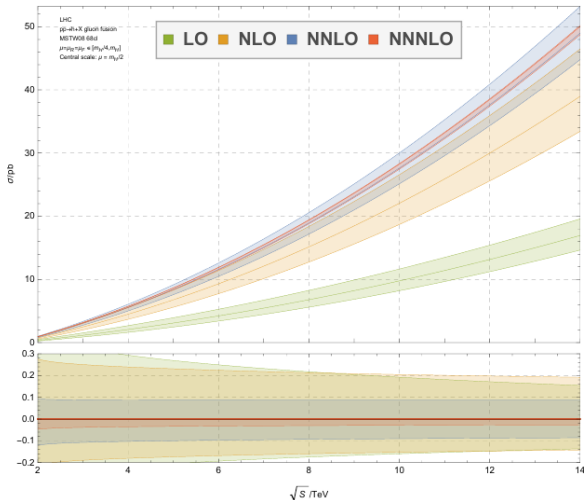
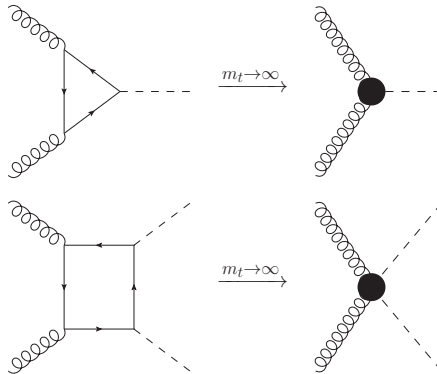


Figure: [Anastasiou et. al 2015]

# Wilson coefficients for Higgs-Boson production



# Wilson coefficients for Higgs-Boson production

- In the limit  $m_t \rightarrow \infty$  the coupling of Higgs-bosons to gluons is given by

$$\mathcal{L}_{\text{eff}} = -C_H \frac{H}{v} \mathcal{O}_1 + \frac{C_{HH}}{2} \left( \frac{H}{v} \right)^2 \mathcal{O}_1$$
$$\mathcal{O}_1 = G_{\mu\nu}^a G^{a\mu\nu}$$

- Inclusive Higgs-boson production known up to N<sup>3</sup>LO  
→  $C_H$  needed at 4 loops
- $C_{HH}$  to 4 loops → N<sup>3</sup>LO Higgs pair-production

# Interlude: Running & Decoupling in QCD

QCD typically renormalized in  $\overline{\text{MS}}$ -scheme  $\rightarrow \alpha_s(\mu)$

$$\beta_{\alpha_s} = \mu^2 \frac{d}{d\mu^2} \frac{\alpha_s}{\pi} \qquad \beta_{\alpha_s} = - \sum_{i \geq 0} \beta_i \left( \frac{\alpha_s}{\pi} \right)^{i+2}$$

$$\beta_0 = \frac{1}{4} \left( 11 - \frac{2}{3} n_f \right)$$

[Gross, Wilczek 1973], [Politzer 1973]

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$\beta_4$  (5-loop) recently calculated:

[Baikov, Chetyrkin, Kühn 2016],

[Herzog, Ruijl, Ueda, Vermaseren 2017],

[Luthe, Maier, Marquard, Schroder 2017]

# QCD beta function

QCD typically renormalized in  $\overline{\text{MS}}$ -scheme  $\rightarrow \alpha_s(\mu)$

$$\beta_{\alpha_s} = \mu^2 \frac{d}{d\mu^2} \frac{\alpha_s}{\pi}$$

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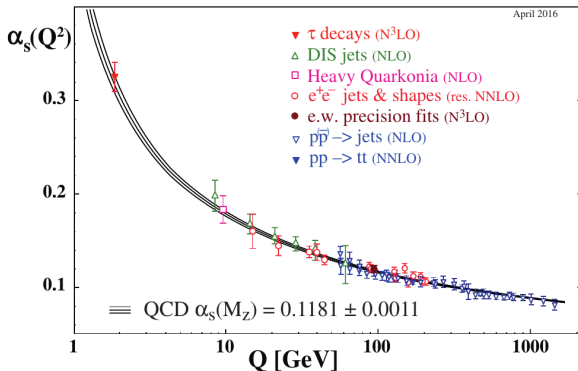


Figure: [PDG 2018]



$\sqrt{s} \ll m_t \rightarrow$  work in effective 5-flavour theory  
 $\rightarrow$  decouple top-quark from running of  $\alpha_s$

$\overline{\text{MS}}$  scheme:

$$\alpha_s^{(5)}(\mu_{\text{dec}}) = \zeta_{\alpha_s} \left( \ln \left( \frac{\mu_{\text{dec}}^2}{m_t^2} \right) \right) \alpha_s^{(6)}(\mu_{\text{dec}})$$

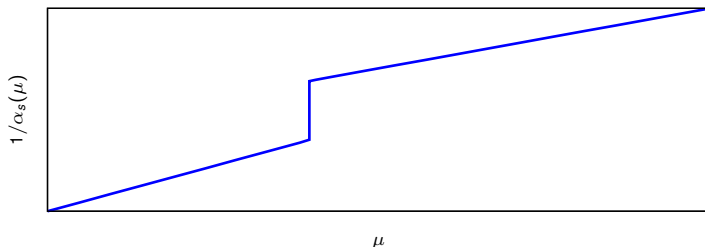
$\zeta_{\alpha_s}$  known to 4 loops for  $N_C = 3$ :

[Chetyrkin, Kühn, Sturm 2015], [Schroder, Steinhauser 2015]

# Decoupling of heavy quarks

$\sqrt{s} \ll m_t \rightarrow$  work in effective 5-flavour theory  
 $\rightarrow$  decouple top-quark from running of  $\alpha_s$

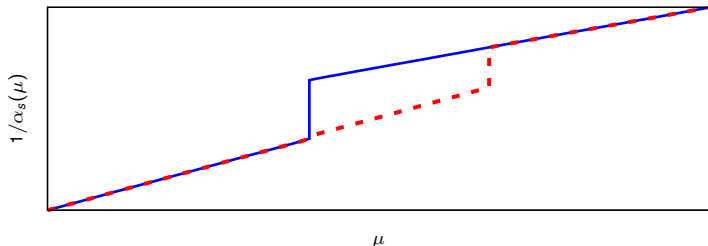
$\rightarrow \alpha_s(\mu)$  is not continuous:



# Decoupling of heavy quarks

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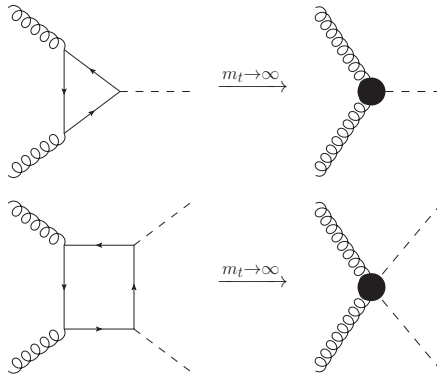


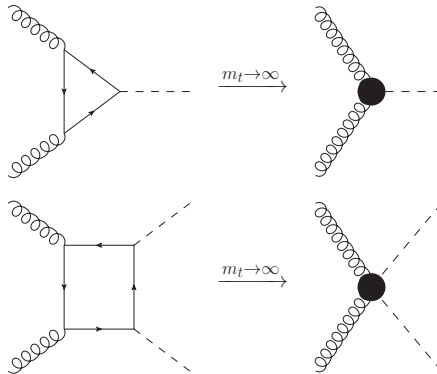
- QCD running & decoupling
- Relations between various mass-schemes
- highest perturbative accuracy available
- Mathematica and C++

RunDec: [Chetyrkin, Kühn, Steinhauser, 2000]

CRunDec: [Schmidt, Steinhauser, 2012]

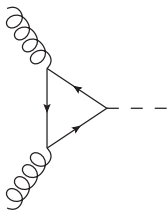
# Decoupling & Higgs-physics



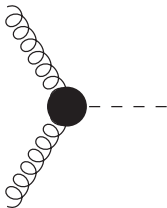


$$\mathcal{L}_{\text{eff}} = -C_H \frac{H}{v} \mathcal{O}_1 + \frac{C_{HH}}{2} \left( \frac{H}{v} \right)^2 \mathcal{O}_1$$

$$\mathcal{O}_1 = G_{\mu\nu}^a G^{a\mu\nu}$$



$$= \epsilon_1^\mu \epsilon_2^\mu (g_{\mu\nu} (q_1 \cdot q_2) - q_{1\nu} q_{2\mu}) F(m_t, m_H)$$

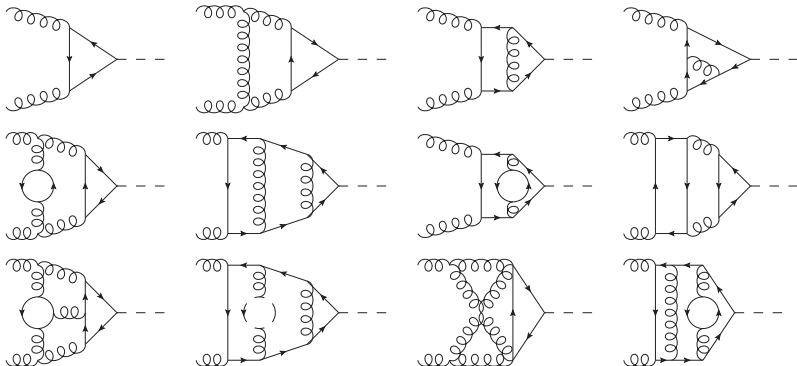


$$= \epsilon_1^\mu \epsilon_2^\mu (g_{\mu\nu} (q_1 \cdot q_2) - q_{1\nu} q_{2\mu}) \frac{C_H}{v}$$

→ At leading order:  $C_H = vF|_{m_t \rightarrow \infty}$



23251 diagrams:



- Expand  $F$  in  $\frac{1}{m_t}$
- technical challenge: "massive tadpoles" up to 4 loops

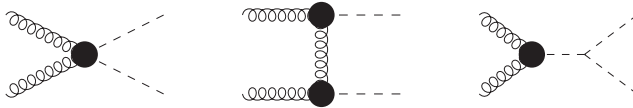
$$C_H = -\frac{\alpha_s^{(5)}}{3\pi} \left\{ 1 + \dots + \left( \frac{\alpha_s^{(5)}}{\pi} \right)^3 \left[ C_A^3 \left( -\frac{110041}{124416} + \frac{1577}{9216} \zeta(3) - \frac{1993}{3456} \ln \frac{\mu^2}{m_t^2} - \frac{77}{576} \ln^2 \frac{\mu^2}{m_t^2} \right) + \dots \right] \right\}$$

- Another way to compute  $C_H$  [Chetyrkin, Kniehl, Steinhauser 1997]:
- Low-Energy-Theorem (LET):

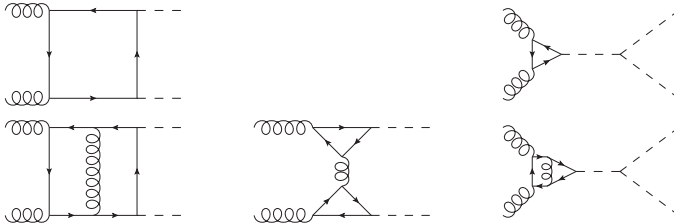
$$C_H = -\frac{2}{m_t^2 \zeta_{\alpha_s}} \frac{\partial}{\partial m_t^2} \zeta_{\alpha_s}$$

- Explicit calculation of  $C_H$  agrees with LET
- Good warmup for computing  $C_{HH}$

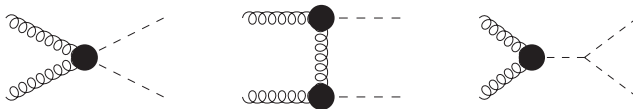
$$C_{HH} \mathcal{A}_{1PI}^{\text{eff}} + C_H^2 \mathcal{A}_{1PR, \lambda=0}^{\text{eff}} + C_H \mathcal{A}_{1PR, \lambda \neq 0}^{\text{eff}}$$



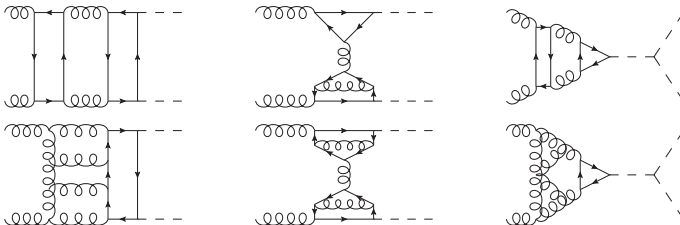
$$= \frac{1}{\zeta_0^3} \left( \mathcal{A}_{1PI}^{\text{full}} + \mathcal{A}_{1PR, \lambda=0}^{\text{full}} + \mathcal{A}_{1PR, \lambda \neq 0}^{\text{full}} \right)$$



$$C_{HH} \mathcal{A}_{1PI}^{\text{eff}} + C_H^2 \mathcal{A}_{1PR, \lambda=0}^{\text{eff}} + C_H \mathcal{A}_{1PR, \lambda \neq 0}^{\text{eff}}$$



$$= \frac{1}{\zeta_3^0} \left( \mathcal{A}_{1PI}^{\text{full}} + \mathcal{A}_{1PR, \lambda=0}^{\text{full}} + \mathcal{A}_{1PR, \lambda \neq 0}^{\text{full}} \right)$$



- 145942 diagrams
- 1 and 2 loop  $C_{HH} = C_H$
- 3 loops:

$$C_{HH} - C_H = \frac{7}{8}C_A^2 - \frac{5}{6}C_A T_F - \frac{11}{8}C_A C_F + \frac{1}{2}C_F T_F + C_F T_F n_l$$

[Grigo, Melnikov, Steinhauser 2014]

- At 4 loops there are 19 colour factors
- 16 of them completed
- 3 work in progress

$$C_{HH} - C_H = \left( \frac{\alpha_s^{(5)}}{\pi} \right)^4 \left[ C_A^3 \left( -\frac{1993}{1728} - \frac{77}{144} \ln \frac{\mu^2}{m_t^2} \right) + \dots \right]$$

- [Spira 2016]:

$$C_{HH} = \frac{1}{4v^2} \left[ \left( \frac{m_t}{\zeta_{\alpha_s}} \frac{\partial}{\partial m_t} \zeta_{\alpha_s} \right)^2 - \frac{m_t^2}{2\zeta_{\alpha_s}} \frac{\partial^2}{\partial m_t^2} \zeta_{\alpha_s} \right]$$

- The 16 colour factors of  $C_{HH}$  we explicitly computed agree

$$\mathcal{L}_{\text{eff}} = -C_H \frac{H}{v} \mathcal{O}_1 + \frac{C_{HH}}{2} \left( \frac{H}{v} \right)^2 \mathcal{O}_1$$

- Explicit calculation of  $C_H$  and  $C_{HH}$  at 4 loops  
→ ingredient for  $N^3\text{LO}$  production cross-sections
- Agreement with LET
- Byproduct: all decoupling constants in QCD at 4 loops